Supplementary Material Available. Tables containing the experimental $1 / T_{1}^{0}, 1 / T_{1.0 b s d}$, and $1 / T_{1 P}$ data corresponding to Figure 4 as well as the line width data corresponding to Figure 5 and the chemical shift and contact shift data corresponding to Figure 6 will appear following these pages in the microfilm edition of this volume of the journal. Photocopies of the supplementary material from this paper only or microfiche ( $105 \times 148 \mathrm{~mm}, 24 \times$ reduction. negatives) containing all of the supplementary material for the papers in this issue may be obtained from the Books and Journals Division, American Chemical Society, 1155 16th St., N.W., Washington, D.C. 20036. Remit check or money order for $\$ 4.00$ for photocopy or $\$ 2.50$ for microfiche, referring to code number JACS-75-6962.

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# Ab Initio Calculation Including Electron Correlation of the Structure and Binding Energy of $\mathrm{BH}_{5}$ and $\mathrm{B}_{2} \mathrm{H}_{7}{ }^{-}$ 

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#### Abstract

Quantum chemical ab initio calculations with and without the inclusion of electron correlation are performed for five possible structures of $\mathrm{BH}_{5}$ and four possible structures of $\mathrm{B}_{2} \mathrm{H}_{7}{ }^{-}$. In the SCF approximation $\mathrm{BH}_{5}$ is not stable with respect to $\mathrm{BH}_{3}+\mathrm{H}_{2}$; with correlation it has a binding energy of $-2 \mathrm{kcal} / \mathrm{mol}$ in a $C_{s}$ geometry. Isomerization is possible via the $C_{2 v}$ structure, the energy of which is $9 \mathrm{kcal} / \mathrm{mol}$ higher. For $\mathrm{B}_{2} \mathrm{H}_{7}^{-}$previous $S C F$ results are confirmed and with correlation the experimental binding energy (with respect to $\mathrm{BH}_{3}+\mathrm{BH}_{4}{ }^{-}$) is reproduced. The ion is symmetric with one linear B-H-B bond.


## I. Introduction

The molecule $\mathrm{BH}_{5}$ is isoelectronic to the ion $\mathrm{CH}_{5}{ }^{+}$that has been studied both experimentally ${ }^{1-3}$ and theoretically. ${ }^{4-8}$ Since $\mathrm{CH}_{5}{ }^{+}$is quite stable as an "ion in space" (its binding energy with respect to $\mathrm{CH}_{4}+\mathrm{H}^{+}$is $-125 \mathrm{kcal} / \mathrm{mol}$ and with respect to $\left.\mathrm{CH}_{3}{ }^{+}+\mathrm{H}_{2}-40 \mathrm{kcal} / \mathrm{mol}\right)$, ${ }^{6-8}$ one could guess that $\mathrm{BH}_{5}$ should be a stable species as well. Recent experimental studies of the protolysis or hydrolysis ${ }^{9-11}$ of $\mathrm{BH}_{4}{ }^{-}$led to the postulate of $\mathrm{BH}_{5}$ as an intermediate in these reactions. A CNDO/2 calculation ${ }^{9}$ of four possible structures of $\mathrm{BH}_{5}$ (belonging to the symmetry groups $D_{3 h}$. $C_{3 v}, C_{4 v}$, and $C_{s}$ ) led to the prediction that $\mathrm{BH}_{5}$ should have $C_{s}$ symmetry (like $\mathrm{CH}_{5}{ }^{+}$) with a binding energy of
$-1.531 \mathrm{au}(\sim-40 \mathrm{eV}$ or $-900 \mathrm{kcal} / \mathrm{mol})$ with respect to B +5 H . Since this value of the binding energy is rather unrealistic (something similar was observed for other boron hydrides like $\mathrm{B}_{2} \mathrm{H}_{6}{ }^{12}$ ) the result from the CNDO calculation could not be regarded as definite and we therefore performed a refined $a b$ initio calculation, including electron correlation, with the CEPA-PNO and PNO-CI methods described elsewhere. ${ }^{14-19}$

The ion $\mathrm{B}_{2} \mathrm{H}_{7}{ }^{-}$that is well-known from experiment ${ }^{20-23}$ is interesting insofar as a structure with a single linear symmetric $\mathrm{B}-\mathrm{H}-\mathrm{B}$ bond has been proposed for this ion. Ab initio studies including correlation of $\mathrm{B}_{2} \mathrm{H}_{6}{ }^{13,24,25}$ that has two $\mathrm{B}-\mathrm{H}-\mathrm{B}$ bonds, of $\mathrm{BeBH}_{5}{ }^{26}$ that is held together by three $\mathrm{Be}-\mathrm{H}-\mathrm{B}$ bonds, and $\mathrm{BeB}_{2} \mathrm{H}_{8}{ }^{26}$ where structures with two

Table I. Optimum Geometry and Binding Energy of $\mathrm{BH}_{5}$ in Different Structures in SCF Approximation ${ }^{a}$

| Con- <br> figuration | Symmetry <br> group | $R_{\mathrm{B}-\mathrm{H}, \mathrm{a}_{0}}$ | $r_{\mathrm{B}-\mathrm{H}, \mathrm{a}_{0}}$ | $\sigma_{\mathrm{B}-\mathrm{H}, \mathrm{a}_{0}}$ | $\alpha, \mathrm{deg}$ | $\beta, \mathrm{deg}$ | $\gamma, \mathrm{deg}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |

${ }^{a}$ The structures and the geometrical parameters are illustrated in Figure 1. The values for configuration $I$ are in parentheses because they do not correspond to a (symmetry-restricted) minimum in the proper sense (see text). ${ }^{b} E_{\mathrm{BIND}}=E_{\mathrm{SCF}}\left(\mathrm{BH}_{5}\right)-E_{\mathrm{SCF}}\left(\mathrm{BH}_{3}\right)-E_{\mathrm{SCF}}\left(\mathrm{H}_{2}\right)$.
and three hydrogen bonds compete have shown that the correlation energy has a marked influence on the binding energy in those systems. We therefore regard a calculation of $\mathrm{B}_{2} \mathrm{H}_{7}{ }^{-}$on the same level of approximation as worthwhile, even though good SCF results have already been published ${ }^{25}$ We also decided to compare competitive structures for $\mathrm{B}_{2} \mathrm{H}_{7}{ }^{-}$with one, two, and three hydrogen bonds, respectively.

## II. Method and Basis Set

We used the IEPA-PNO, PNO-CI, and CEPA-PNO methods ${ }^{14-19}$ that were recently described in detail ${ }^{18}$ and applied to first row ${ }^{27}$ and second row hydrides ${ }^{28}$ as well as to other molecules. ${ }^{29}$ We always start with an SCF calculation and then transform the SCF-MO's to localized ones according to Boys, ${ }^{30}$ before the calculation of the correlation effects is performed.

In the IEPA (independent electron pair approximation) scheme the correlation energy of the different electron pairs is computed independently, in CEPA (coupled electron pair approximation) the interaction of the electron pairs is taken into account, and Cl means configuration interaction in the usual sense, limited to double substitutions. In all three schemes the pair correlation functions are calculated in a PNO (pair natural orbital) expansion. The CEPA values are usually closest to the exact ones. ${ }^{17.27-29}$ The basis set consists of Gaussian lobes from which $p$ and $d$ functions are constructed as described in ref. 31. A "small basis" 33 was used for the geometry optimizations. It consists of the ( $8 \mathrm{~s}, 4 \mathrm{p}$ ) Huzinaga basis ${ }^{32}$ in the contraction ( $5,1,1,1 ; 3,1$ ) for $B$ and the $4 s$ basis in the contraction $(3,1)$ for $H$.

A "large basis" ${ }^{33}$ was taken for the final calculations at the computed equilibrium geometries. It differs from the "small basis" by the addition of polarization functions, namely a d set $(\eta=0.61)$ for $B$ and a $p$ set $(\eta=0.61)$ for $H$. The $p$ basis of $B$ was tentatively used in the contraction ( $2,1,1$ ), but pilot calculations indicated, in agreement with the findings of Ahlrichs, ${ }^{13}$ that the p function with the smallest $\eta$ could be left out without affecting the results for the binding energies and the energy differences between the isomers of $\mathrm{BH}_{5}$.

The computer times with the "large basis" were by about a factor 5-10 larger than those with the "small basis." One SCF point for $\mathrm{BH}_{5}$ needed $\sim 3.5 \mathrm{~min}$ with the small and $\sim 32 \mathrm{~min}$ with the large basis, one point including correlation needed 82 min . For one SCF point of $\mathrm{B}_{2} \mathrm{H}_{7}^{-}$in the small basis 50 min and in the large basis 220 min were necessary. The calculation of correlation with the large basis is in principle possible for $\mathrm{B}_{2} \mathrm{H}_{7}$, but we chose a less time-consuming compromise as described in section VIII that needed only $2 \times 150 \mathrm{~min}$. These values of computer time are for the Telefunken TR 440.

## III. Search for the Equilibrium Structure of $\mathrm{BH}_{5}$ in the SCF Approximation

Five types of structures as illustrated in Figure 1 were considered. They are the same as in the study of $\mathrm{CH}_{5}{ }^{+6.8}$




conf. $\boldsymbol{Y}$


Figure 1. Five possible structures of $\mathrm{BH}_{5}$.

Starting values for the geometrical parameters were those of $\mathrm{BH}_{3} .{ }^{34}$ The following structural parameters were varied (see Figure 1): (configuration I) $r_{\mathrm{BH}}, R_{\mathrm{BH}}, \alpha, \beta$ (for the $\mathrm{BH}_{3}$ subunit local $C_{3 v}$ symmetry was retained); (configuration II) $R_{\mathrm{BH}}, R_{\mathrm{HH}}, \alpha$; (configuration III) $r_{\mathrm{BH}}, R_{\mathrm{BH}}$; (configuration IV) $r_{\mathrm{BH}}, R_{\mathrm{BH}}, \alpha$; (configuration V) $r_{\mathrm{BH}}, \sigma_{\mathrm{BH}}$, $R_{\mathrm{BH}}, \alpha$.

The results of the geometry optimization for the different structures are collected in Table 1. The energies given in Table I were calculated with the "large basis", but the "small basis" was used for the geometry optimization. For structure II there was a continuous increase of the energy on approach of the $\mathrm{H}_{2}$ to $\mathrm{BH}_{3}$ and no minimum could be found by relaxation of the other geometrical parameters. The same was true for structure I, but here one could force the energy to go through a minimum if one constrained the out-of-plane angle $\alpha$ of the $\mathrm{BH}_{3}$ subunit to be constant and to be larger than $5^{\circ}$. The energy curve for $\alpha=8^{\circ}$ as a function of the distance $R$ between the B atom and the midpoint of the $\mathrm{H}_{2}$ subunit is illustrated in Figure 2. The relative


Figure 2. SCF energy of $\mathrm{BH}_{5}$ in configuration I for $\alpha$ and $r_{\mathrm{H}-\mathrm{H}}$ fixed.
minimum (i.e., the binding energy of $\mathrm{BH}_{5}$ with respect to $\mathrm{H}_{2}$ and bent $\mathrm{BH}_{3}$ ) increases with $\alpha$ (for $5^{\circ}<\alpha<30^{\circ}$ ) which means that pyramidal $\mathrm{BH}_{3}$ with an sp ${ }^{3}$-type unoccupied AO is better able to bind $\mathrm{H}_{2}$ via a three-center twoelectron bond than is planar $\mathrm{BH}_{3}$ with a p-type unoccupied AO. However, the energy required to deform $\mathrm{BH}_{3}$ by the out-of-plane angle $\alpha$ increases faster with $\alpha$ so that the absolutely lowest of the constrained minima occurs for the smallest possible value of $\alpha$, i.e., for $\alpha=5^{\circ}$. The energy corresponding to this constrained minimum is, as can be seen from Table I, only $1.4 \mathrm{kcal} / \mathrm{mol}$ above the sum of the energies of $\mathrm{BH}_{3}$ and $\mathrm{H}_{2}$. For structure II not even a constrained minimum of this kind could be found. For the structures III, IV, and $V$ one does find minima but they are energetically much above $\mathrm{BH}_{3}$ and $\mathrm{H}_{2}$. Note that all values for $E_{\text {bind }}$ in Table I are positive, which means that $\mathrm{BH}_{5}$ has a higher energy than $\mathrm{BH}_{3}+\mathrm{H}_{2}$. It is unlikely that for any of these structures there exists a real local minimum; we rather expect that the energy goes down in any case if one relaxes the symmetry restrictions. We conclude that in the SCF approximation $\mathrm{BH}_{5}$ is not bound with respect to $\mathrm{BH}_{3}$ $+\mathrm{H}_{2}$.

All the energy differences quoted in this paper refer to the minima of the potential curves and not to the zero-point vibrational levels.

## IV. The Correlation Energy of $\mathrm{BH}_{5}$

We have computed the correlation energy for the structures I $\left(C_{s}\right)$ and $\mathrm{V}\left(C_{2 v}\right)$ and, of course, of $\mathrm{BH}_{3}$ and $\mathrm{H}_{2}$ in the same basis. There was no reason to expect that structures III or IV would become competitive if electron correlation is taken care of. Inclusion of correlation in the CEPA scheme stabilizes the $C_{2 v}$ structure by $12.4 \mathrm{kcal} / \mathrm{mol}$ with respect to $\mathrm{BH}_{3}+\mathrm{H}_{2}$ so that with inclusion of correlation $\mathrm{BH}_{5}$ in $C_{2 v}$ symmetry lies only $7 \mathrm{kcal} / \mathrm{mol}$ above the energy of the separated systems $\mathrm{BH}_{3}+\mathrm{H}_{2}$. We did not reoptimize the geometry (obtained from SCF calculations) of this $C_{2 v}$ structure. Pilot calculations indicate that such a reoptimization can only lower the energy by 1 or at most $2 \mathrm{kcal} / \mathrm{mol}$.
The $C_{s}$ structure becomes stable with respect to $\mathrm{BH}_{3}+$ $\mathrm{H}_{2}$ if electron correlation is included. Since there was no minimum of the SCF potential surface in $C_{s}$ geometry we had to optimize the geometry with correlation (see section 1II). The equilibrium structure was found for the following values (see Figure 1): $R_{\mathrm{B}-\mathrm{H}}=2.90 \mathrm{a}_{0}, r_{\mathrm{B}-\mathrm{H}}=2.25 \mathrm{a}_{0}, r_{\mathrm{H}-\mathrm{H}}$ $=1.48 \mathrm{a}_{0}, \alpha=8^{\circ}, \beta=29.6^{\circ}$. In Figure 3 the total energy as well as the correlation contribution to the binding energy is plotted as a function of the distance $R$ between B and the midpoint of $\mathrm{H}_{2}$. At the equilibrium geometry the binding energy is $\sim-2 \mathrm{kcal} / \mathrm{mol}$, the correlation contribution is


Figure 3. Correlation energy and total energy of $\mathrm{BH}_{5}$ in configuration I for $\alpha$ and $r_{\mathrm{HH}}$ fixed.

Table II. Total Energies in Different Approximations (in au) for Optimized Structures (Large Basis)

| Molecule | $-E_{\text {SCF }}$ | $-E_{\text {IEPA }}$ | $-E_{\text {CEPA }}$ | $-E_{\text {PNOCI }}$ |
| :--- | :---: | ---: | ---: | ---: |
| $\mathrm{H}_{2}$ | 1.129086 | 1.163913 | 1.163913 | 1.163913 |
| $\mathrm{BH}_{3}$ | 26.387012 | 26.512176 | 26.503314 | 26.498504 |
| $\mathrm{BH}_{5}\left[C_{s}\right]$ | $(27.507790)^{a}$ | 27.683820 | 27.669950 | 27.660190 |
| $\mathrm{BH}_{s}\left[C_{2 v}\right]$ | $27.513797{ }^{b}$ |  |  |  |

$a_{\text {For the }}$ minimum of the CEPA calculation. ${ }^{b}$ For the constrained SCF minimum (see text).
$\sim-7 \mathrm{kcal} / \mathrm{mol}$, and the SCF repulsion $\sim+5 \mathrm{kcal} / \mathrm{mol}$. The pertinent total energies for the equilibrium geometries are collected in Table II.

Obviously the fact that $\mathrm{BH}_{5}$ is bound with respect to $\mathrm{BH}_{3}$ $+\mathrm{H}_{2}$ is exclusively due to correlation. Before we comment on this somewhat unusual case of a chemical bond we want to say a few words about the individual contributions to the correlation energy in $\mathrm{BH}_{5}, \mathrm{BH}_{3}$, and $\mathrm{H}_{2}$.

The correlation energy of a molecule computed in the coupled electron pair approximation (CEPA) can be decomposed into independent electron pair (IEPA) contributions $\epsilon_{\mu}$ and coupling correlation terms $\Delta \epsilon_{\mu \nu}$, where $\mu, \nu$ count the pairs. The bulk of the change in correlation energy is in the IEPA contributions $\epsilon_{\mu}$ which are collected in Table III. The localized MO's in $\mathrm{BH}_{5}$ can directly be compared with those of $\mathrm{BH}_{3}$ and $\mathrm{H}_{2}$ since $\mathrm{BH}_{5}$ in structure 1 contains rather distinct $\mathrm{BH}_{3}$ and $\mathrm{H}_{2}$ subunits. One sees from Table III that the contributions of the $\mathrm{BH}_{3}$ subunit in $\mathrm{BH}_{5}$ to the correlation energy are smaller in absolute value than in isolated $\mathrm{BH}_{3}$ by $0.12480-0.12270=0.0021$ au but that the correlation energy of the $\mathrm{H}_{2}$ subunit is larger in $\mathrm{BH}_{5}$ by 0.00182 in absolute value so that the intrasubsystem correlation energy changes by only 0.0004 au . The intersubsystem correlation energy amounts to $0.01828 \approx 13$ $\mathrm{kcal} / \mathrm{mol}$ and equals practically the total change in the IEPA correlation energy. The pair coupling terms reduce this difference to the CEPA value of $\sim 7 \mathrm{kcal} / \mathrm{mol}$.

From the experience with other systems we conclude that the CEPA scheme predicts energy differences with chemical accuracy, i.e., with an error of the order of magnitude of $1 \mathrm{kcal} / \mathrm{mol}$ or possibly less. In the present case refinement of the calculation is likely to increase the value of the binding energy.

## V. Analysis of the Binding in the $\mathrm{BH}_{5}$

In view of the small binding energy of only $-2 \mathrm{kcal} / \mathrm{mol}$ between $\mathrm{BH}_{3}$ and $\mathrm{H}_{2}$ one wonders whether $\mathrm{BH}_{5}$ should be

Table III. Pair Contributions (Localized) to the Correlation
Energy (Large Basis) (in au) for $\mathrm{BH}_{5}, \mathrm{BH}_{3}$, and $\mathrm{H}_{2}$ in
Their Equilibrium Geometries

| Molecule | Subunit | $\underset{\text { Pair }}{\text { symbol }}{ }^{a}$ | No. of equivalent pairs | Pair contribution | Sum |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $\begin{aligned} & \mathrm{H}_{2} \\ & \mathrm{BH}_{3} \end{aligned}$ |  | $-\epsilon_{h}$ | 1 | 0.03483 | 0.03483 |
|  |  | $-\epsilon_{b}$ | 3 | 0.03129 | 0.09386 |
|  |  | $-\epsilon_{\mathrm{bb}}{ }^{\prime}$ | 3 | 0.01032 | 0.03094 |
|  |  |  |  |  | 0.12480 |
| $\begin{aligned} & \mathrm{BH}_{5} \\ & \mathrm{BH}_{5} \end{aligned}$ | $\begin{aligned} & \mathrm{H}_{2} \\ & \mathrm{BH}_{3} \end{aligned}$ | $-\epsilon_{h}$ | 1 | 0.03665 | 0.03665 |
|  |  | $-\epsilon_{\mathrm{b}}$ | 1 | 0.03077 | 0.03077 |
|  |  | $-\epsilon_{\mathrm{b}}$, | 2 | 0.03055 | 0.06110 |
|  |  | $-\epsilon_{\mathrm{b}}{ }^{\prime}$ | 2 | 0.01020 | 0.02040 |
|  |  | $-\epsilon_{\mathrm{b}} \mathrm{b}^{\prime \prime}$ | 1 | 0.01043 | 0.01043 |
|  |  |  |  |  | 0.12270 |
| $\mathrm{BH}_{s}$ | Inter$\mathrm{H}_{2} / \mathrm{BH}_{3}$ | $-\epsilon_{\mathrm{bh}}$ | 1 | 0.00550 | 0.00550 |
|  |  | $-\epsilon_{\mathrm{b}}{ }^{\prime} \mathrm{h}$ | 2 | 0.00639 | 0.01278 |
|  |  |  |  |  | 0.01828 |

$a^{\text {Subscript }} \mathrm{h}$ refers to a $\mathrm{H}-\mathrm{H}$ bond and b to a BH bond.

Table IV. Mulliken Gross Charges and Overlap Populations for the Localized MO's in $\mathrm{BH}_{5}$

|  | B | $\mathrm{H}(1)$ | $\mathrm{H}(2)$ | $\mathrm{H}(3)$ | $\mathrm{H}(4)$ | $\mathrm{H}(5)$ |
| :--- | ---: | :---: | :---: | :---: | :---: | :---: |
|  | 0.77 | $\mathrm{BH}(2)$ | Two-Center Bond |  |  |  |
| B | 0.05 | -0.03 |  |  |  |  |
| $\mathrm{H}(1)$ | -0.03 |  |  |  |  |  |
| $\mathrm{H}(2)$ | 0.86 | -0.07 | 1.27 |  |  |  |
| $\mathrm{H}(3)$ | -0.04 | 0.01 | -0.06 | -0.02 |  |  |
| $\mathrm{H}(4)$ | -0.06 | 0.02 | -0.08 | 0.02 | -0.02 |  |
| $\mathrm{H}(5)$ | 0.04 | -0.01 | 0.05 | 0.00 | -0.02 | 0.03 |
|  |  | $\mathrm{BH}(5) \mathrm{H}(6)$ Three-Center Bond |  |  |  |  |
| B | 0.18 |  |  |  |  |  |
| $\mathrm{H}(1)$ | 0 | -0.01 |  |  |  |  |
| $\mathrm{H}(2)$ | 0 | 0 | -0.01 |  |  |  |
| $\mathrm{H}(3)$ | 0 | 0 | 0 | -0.01 |  |  |
| $\mathrm{H}(4)$ | 0.12 | -0.02 | -0.02 | -0.03 | 0.93 |  |
| $\mathrm{H}(5)$ | 0.12 | -0.02 | -0.03 | -0.02 | $0.74{ }^{a}$ | 0.93 |

${ }^{a}$ For comparison: the overlap population of isolated $\mathrm{H}_{2}$ in the same basis is 0.83 .
regarded as a van der Waals complex rather than a chemically bound molecule. Like in genuine van der Waals complexes the intersystem correlation energy (dispersion) is mainly responsible for the existence of a stable $\mathrm{BH}_{5}$. We nevertheless prefer to look at $\mathrm{BH}_{5}$ as a chemically bound molecule and this for the following reasons: (1) A van der Waals bond between $\mathrm{H}_{2}$ and $\mathrm{BH}_{3}$ should have a still much smaller binding energy (the well depth of the van der Waals minimum of the system $\mathrm{He}-\mathrm{H}_{2}$ is, e.g., $0.04 \mathrm{kcal} / \mathrm{mol} ;{ }^{35}$ other typical van der Waals minima are of the same order of magnitude). (2) The distance between the two subunits is much smaller than typical van der Waals distances (the distance between He and the midpoint of $\mathrm{H}_{2}$ at the van der Waals minimum is $6.3 \mathrm{a}_{0}{ }^{35}$ ). (3) The $\mathrm{BH}_{3}$ and $\mathrm{H}_{2}$ subunits in $\mathrm{BH}_{5}$ are appreciably deformed compared with the isolated $\mathrm{BH}_{3}$ and $\mathrm{H}_{2}$. The out-of-plane angle in the $\mathrm{BH}_{3}$ subunit is $8^{\circ}$ and the bond distance of the $\mathrm{H}_{2}$ subunit is $1.48 \mathrm{a}_{0}$ compared with $1.40 \mathrm{a}_{0}$ in isolated $\mathrm{H}_{2}$.

The localized MO (LMO) that represents the three-center BHH bond is plotted in Figure 4. One sees from Figure 4 and also from the population analysis in Table IV that this LMO is mainly localized in the $\mathrm{H}-\mathrm{H}$ bond, but that there is a nonnegligible charge transfer ( 0.2 electron) to the B atom and a significant $\mathrm{B}-\mathrm{H}$ overlap population. The LMO's that represent ordinary B-H bonds are very much like those in $\mathrm{BH}_{3}$.

We may compare $\mathrm{BH}_{5}$ with other molecules which contain a $\mathrm{H}_{2}$ subunit that form a three-center two-electron


Figure 4. Contour line diagram in the BHH plane of the localized MO of $\mathrm{BH}_{5}$ (equilibrium structure) that represents the three-center BHH bond.

Table V. Compounds in $C_{s}$ Symmetry with Three-Center Two Electron Bonds ${ }^{a}$

|  | $R, \mathrm{a}_{0}$ | $r_{\mathrm{H}-\mathrm{H}}, \mathrm{a}_{0}$ | $E_{\mathrm{BIND}}$, <br> $\mathrm{kcal} / \mathrm{mol}$ | $\Delta E_{\mathrm{CORR}}$, <br> $\mathrm{kcal} / \mathrm{mol}$ |
| :--- | :---: | :---: | :---: | :---: |
| $\mathrm{LiH}_{2}{ }^{+}$ | 3.75 | 1.41 | -5.1 | 0 |
| $\mathrm{BeH}_{3}{ }^{+}$ | 3.04 | 1.47 | -24.7 | 0 |
| $\mathrm{BH}_{3}{ }^{+}$ | 2.80 | 1.48 | -2.0 | 7.0 |
| $\mathrm{CH}_{5}{ }^{+}$ | 2.24 | 1.67 | -43.0 | 21.0 |
| $\mathrm{H}_{3}{ }^{+}$ | 1.43 | 1.65 | -107.1 | 4.0 |

$a_{\mathrm{K}}=\mathrm{Li}, \mathrm{Be}, \mathrm{B}, \mathrm{C} ; R=$ distance of nucleus K from the center of mass of $\mathrm{H}_{2} ; r_{\mathrm{H}-\mathrm{H}}=$ distance between the H atoms in the $\mathrm{H}_{2}$ subunit; $E_{\mathrm{BIND}}=E_{\mathrm{KH} n}{ }^{\text {total }}-E_{\mathrm{KH} n-2}{ }^{\text {total }}-E_{\mathrm{H}_{2}}^{\text {total } ; \Delta E_{\mathrm{CORR}}}=$ correlation contribution to the binding energy.
bond with an electron deficient subunit namely with $\mathrm{LiH}_{2}{ }^{+},{ }^{36} \mathrm{BeH}_{3}{ }^{+},{ }^{37} \mathrm{CH}_{5}+,{ }^{6-8}$ and $\mathrm{H}_{3}+{ }^{+38-40}$ In Table V the bond distances and binding energies of these molecules are compared with those of $\mathrm{BH}_{5}$. As far as the distance between the heavy atom and the $\mathrm{H}_{2}$ subunit is concerned and with respect to the change of the $\mathrm{H}-\mathrm{H}$ distance $\mathrm{BH}_{5}$ fits well in between $\mathrm{BeH}_{3}{ }^{+}$and $\mathrm{CH}_{5}{ }^{+}$. Also the contribution of the change in correlation energy to the binding energy of $\mathrm{BH}_{5}$ is between that of $\mathrm{BeH}_{3}{ }^{+}$and $\mathrm{CH}_{5}{ }^{+}$. Only the binding energy of $\mathrm{BH}_{5}$ does not fit at all. However, all the molecules to which we can compare $\mathrm{BH}_{5}$ are ionic and their binding energies contain an appreciable amount of induction energy (interaction between the charge of, e.g., $\mathrm{BeH}^{+}$and the induced dipole in $\mathrm{H}_{2}$ ), and electrostatic energy (between the charge and the quadrupole moment of $\mathrm{H}_{2}$ ), which are not present in $\mathrm{BH}_{5}$. The best description of binding between $\mathrm{BH}_{3}$ and $\mathrm{H}_{2}$ is probably to say that it is a three-center bond for which $\mathrm{H}_{2}$ is donor and $\mathrm{BH}_{3}$ acceptor and that this bond is effectively rather weak due to the fact that $\mathrm{H}_{2}$ is a rather poor electron donor and to the strong electron repulsion between the subunits.

This electron repulsion is overestimated in the HartreeFock approximation, but taken care of correctly if electron correlation is allowed for. To conclude this section we compare the equilibrium geometry obtained for $\mathrm{BH}_{5}$ by Olah et al. with the CNDO/ 2 approximation. This is done in Table VI. Although the binding energy from CNDO/ 2 of $\sim-900$ $\mathrm{kcal} / \mathrm{mol}$ with respect to $\mathrm{B}+5 \mathrm{H}$ is unrealistic (we get $\omega-380 \mathrm{kcal} / \mathrm{mol}$ ) the equilibrium geometry is not so bad, at least as far as the angle $\alpha$ and the (uninteresting) distance $r_{B H}$ is concerned.





Figure 5. Four possible structures of $\mathrm{B}_{2} \mathrm{H}_{7}^{-}$

Table VI. Equilibrium Values of $\mathrm{BH}_{s}$ in $C_{S}$ Geometry

|  | $\mathrm{CNDO} / 2$ | This work |
| :---: | :---: | :---: |
| $\gamma, \mathrm{deg}$ | 108 | 119.5 |
| $r_{\mathrm{B}-\mathrm{H}}, \mathrm{A}$ | 1.20 | 1.19 |
| $R_{\mathrm{B}-\mathrm{H}}, \AA$ | 1.31 | 1.53 |
| $r_{\mathrm{H}-\mathrm{H}, \mathrm{A}}$ | 0.92 | 0.78 |

Table VII. Equilibrium Values of $\mathrm{B}_{2} \mathrm{H}_{7}{ }^{-}$with $D_{3} h$ Structure

| Structure I | Hall et al. ${ }^{2 \mathrm{~s}}$ | This work |
| :--- | :---: | :---: |
| $r_{\mathrm{B}-\mathrm{H}}, \AA$ | 1.223 | $1.225\left(2.315 \mathrm{a}_{0}\right)$ |
| $R_{\mathrm{B}-\mathrm{H}}, \AA$ | 1.309 | $1.323\left(2.550 \mathrm{a}_{0}\right)$ |
| $\gamma, \mathrm{deg}$ | 113.8 | 114.3 |
| $\alpha, \operatorname{deg}$ | 14.7 | 14.0 |
| $E_{\text {BIND }}(\mathrm{SCF}), \mathrm{kcal} / \mathrm{mol}$ | -25.1 | -24.8 |
| $E_{\mathrm{BIND}}(\mathrm{CEPA}), \mathrm{kcal} / \mathrm{mol}$ |  | -28.1 |

## VI. Comment on the Experiments Concerning $\mathrm{BH}_{5}$

The existence of $\mathrm{BH}_{5}$ as postulated from experiments ${ }^{9-11}$ is confirmed by our calculations. To explain the isotope exchange results ${ }^{9-11}$ it is not sufficient that $\mathrm{BH}_{5}$ exists; one has also to postulate that the barrier for isomerization is sufficiently small. From our calculation this barrier, namely the energy difference between structures I and $V$, equals $\sim 9$ $\mathrm{kcal} / \mathrm{mol}$ (CEPA); however, the barrier is above the dissociation limit to $\mathrm{BH}_{3}+\mathrm{H}_{2}$. One would therefore expect that the decomposition of $\mathrm{BH}_{5}$ into $\mathrm{BH}_{3}+\mathrm{H}_{2}$ is much faster than the isomerization, which is in contradiction to the experimental findings.

A possible explanation of this discrepancy is that the experiments were done in solution whereas the calculations were done for molecules "in space". Both $\mathrm{BH}_{4}{ }^{-}$and the
protolyzing $\mathrm{H}^{+}$are solvated and consequently the $\mathrm{BH}_{5}$ formed is trapped in a solvating sphere. Thus $\mathrm{BH}_{5}$ has probably enough energy to isomerize several times before $\mathrm{BH}_{3}$ and $\mathrm{H}_{2}$ can "escape" from the solvation sphere.

## VII. SCF Calculations of $\mathbf{B}_{2} \mathbf{H}_{7}^{-}$

Four possible structures of $\mathrm{B}_{2} \mathrm{H}_{7}^{-}$are illustrated in Figure 5. Hall, Marynick, and Lipscomb ${ }^{25}$ performed an SCF calculation of structure I of $\mathrm{B}_{2} \mathrm{H}_{7}{ }^{-}$with a linear $\mathrm{B}-\mathrm{H}-\mathrm{B}$ bridge, using a minimal STO basis. They obtained the structure parameters that are compared in Table VII with our results from SCF calculations with the "large" basis. We have also performed SCF calculations for structure II with two (bent) B-H-B bridges and structures III and IV with three bridges. The optimum geometries for these structures are indicated in Table VIII. We have optimized structures II, III, and IV under the contraint that they have $C_{s}$ and $C_{2 c}$ symmetry, respectively, such that the two bridges in structure II and two of the three bridges in structures 111 and IV are symmetrically equivalent. Also the local threefold axes in structures II and IV were imposed. Under these constraints minima for these structures were found. The minimum for structure III lies $250 \mathrm{kcal} / \mathrm{mol}$ and that of structure IV $200 \mathrm{kcal} / \mathrm{mol}$ above the sum of the energies of $\mathrm{BH}_{4}$ and $\mathrm{BH}_{3}$, whereas both structures I and II are bound with respect to $\mathrm{BH}_{3}+\mathrm{BH}_{4}{ }^{-}$, namely structure I by 25 $\mathrm{kcal} / \mathrm{mol}$ and structure II by $6 \mathrm{kcal} / \mathrm{mol}$. This clearly indicates that bonding through a single $\mathrm{B}-\mathrm{H}-\mathrm{B}$ bond is most effective. In structure II one recognizes $\mathrm{BH}_{3}$ and $\mathrm{BH}_{4}^{-}$subunits, whereas in structure I two equivalent $\mathrm{BH}_{3}$ subunits and a $\mathrm{H}^{-}$in the bridge are found. The ion in structure I turns out to be stable with respect to asymmetric deformations, in agreement with Hall et al., 25 i.e., it has $D_{3 h}$ symmetry although only $C_{2 v}$ symmetry was imposed. (We did not investigate the possible difference in energy between a "staggered" and "eclipsed" structure; this difference is probably very small.) While the B-H-B bond in structure 1 prefers to be symmetric, the two $\mathrm{B}-\mathrm{H}-\mathrm{B}$ bonds in structure II are definitely asymmetric. If one forces them to be symmetric the energy is raised by $\sim 60 \mathrm{kcal} / \mathrm{mol}$. We did not investigate whether in going from structure II to structure 1 there is a barrier, but we think that this is unlikely and that structure II does not represent a local minimum. Structures III and IV have unexpectedly high energies and unexpectedly large BH distances in the bridges.

## VIII. Correlation Energy of $\mathrm{B}_{2} \mathrm{H}_{7}{ }^{-}$

For $\mathrm{B}_{2} \mathrm{H}_{7}^{-}$a full calculation with the "large" basis and with inclusion of correlation is possible but somewhat time consuming. We therefore performed the calculation of the correlation energy "in pieces". For structure I we performed a calculation with the large basis for the "left" half of the molecule plus the $\mathrm{B}-\mathrm{H}-\mathrm{B}$ bond and the small basis

Table VIII. Equilibrium Values of $\mathrm{B}_{2} \mathrm{H}_{2}-$ in $C_{S}$ Symmetry and $C_{2 v}$ Symmetry

| Structure II | SCF-optimized values $C_{S}$ symm | Structure III | SCF- <br> optimized values $\mathrm{C}_{2 U} \mathrm{symm}$ | Structure IV | SCFoptimized values $\mathrm{C}_{2 \nu}$ symm |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $r_{\mathrm{B}-\mathrm{H}}, \mathrm{a}_{0}$ | 2.30 |  | 2.30 | $r_{\text {B-H }}, \mathrm{a}_{0}$ | 2.30 |
| $R_{\mathrm{B}-\mathrm{H}}, \mathrm{a}_{\mathrm{o}}$ | 4.00 | $R_{\text {B-H }}, \mathrm{a}_{0}$ | 3.15 | $R_{\mathrm{B}-\mathrm{H}}$, a ${ }_{0}$ | 3.22 |
| $\tau_{\mathrm{B}-\mathrm{H}}$ | 2.30 |  |  | $\tau_{\mathrm{B}-\mathrm{H}}, \mathrm{a}_{0}$ | 2.60 |
| $\alpha ; \beta ; \gamma, \operatorname{deg}$ | 10.0; 57.0; 110.0 | $\alpha$, deg | 65 | $\alpha$, deg | 72 |
| $\delta ; \epsilon$ | 110.0; 112.0 | $\gamma$, deg | 120 | $\gamma$, deg | 120 |
| $E_{\mathrm{BIND}}(\mathrm{SCF}), \mathrm{kcal} / \mathrm{mol}$ | $-6.3$ | $E_{\text {BIND }}(\mathrm{SCF}), \mathrm{kcal} / \mathrm{mol}$ | $+250^{a}$ | $E_{\text {BIND }}(\mathrm{SCF}), \mathrm{kcal} / \mathrm{mol}$ | +196 |
| $E_{\text {BIND }}(\mathrm{CEPA}), \mathrm{kcal} / \mathrm{mol}$ | $-10.0$ |  |  |  |  |

[^0]Table IX. Energies (in au) of $\mathrm{B}_{2} \mathrm{H}_{7}$ in Four Different Structures

| Molecule | Large basis |  |  | $\begin{gathered} \text { Small } \\ \text { basis } \\ -E_{S C F} \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: |
|  | $-E_{\text {SCF }}$ | $-E_{\text {IEPA }}{ }^{a}$ | $-E_{\text {CEPA }}{ }^{\text {a }}$ |  |
| $\mathrm{BH}_{3}$ | 26.3870 | 26.5121 | 26.5033 | 26.3757 |
| $\mathrm{BH}_{4}^{-}$ | 26.9689 | 27.1475 | 27.1297 | 26.9560 |
| $\mathrm{B}_{2} \mathrm{H}_{7}-[\mathrm{I}]$ | 53.3953 | 53.7080 | 53.6778 | 53.3725 |
| $\mathrm{B}_{2} \mathrm{H}_{7}-$ [II] | 53.3659 | 53.6779 | 53.6489 | 53.3435 |
| $\mathrm{B}_{2} \mathrm{H}_{7}{ }^{-}$[III] |  |  |  | 52.9341 |
| $\mathrm{B}_{2} \mathrm{H}_{7}^{-}$[IV] | 53.0331 |  |  | 53.0331 |

for the "right" half. For the left half plus the bridge the full correlation energy in the IEPA-PNO, CEPA-PNO, and PNO-CI schemes was calculated. From twice this correlation energy the contribution of the $\mathrm{B}-\mathrm{H}-\mathrm{B}$ bond was subtracted. In doing so one ignores the interaction terms between BH bonds on different sides of the ion and also some nonadditivity corrections. The error is, as calculations on comparable molecules have shown, at most of the order of 1 $\mathrm{kcal} / \mathrm{mol}$. For structure II where the "left" and the "right" half of the molecule are not equivalent, two calculations had to be performed. The total energies with and without correlation are collected in Table IX. Some test calculations indicated that the geometries are not affected appreciably by correlation, we therefore did not reoptimize the geometries. The binding energies which in SCF approximation were $\sim-25$ and $\sim-6 \mathrm{kcal} / \mathrm{mol}$ or structures I and II, respectively, are found by CEPA to be -28 and $-10 \mathrm{kcal} /$ mol ; the experimental value is $-30 \mathrm{kcal} / \mathrm{mol}$, ${ }^{15}$ so the agreement is rather good. The contribution of the change in correlation energy to the binding energy is not so spectacular as in $\mathrm{B}_{2} \mathrm{H}_{6}$ or in $\mathrm{BH}_{5}$.

In spite of the noticeable binding energy of $\sim-30 \mathrm{kcal} /$ mol of $\mathrm{B}_{2} \mathrm{H}_{7}^{-}$with respect to $\mathrm{BH}_{3}+\mathrm{BH}_{4}^{-}$the $\mathrm{B}-\mathrm{H}-\mathrm{B}$ bond turns out to be very "soft", in the sense that the energy is very insensitive to changes in the B-B distance. Variation of this distance by $\pm 0.1 a_{0}$ raises the energy by only 0.002 au . This is found both in calculations with and without correlation.

## IX. Conclusion

$\mathrm{BH}_{5}$ is stable in a $C_{s}$ structure (analogous to that of $\mathrm{CH}_{5}{ }^{+}$) by as little as $2 \mathrm{kcal} / \mathrm{mol}$ with respect to $\mathrm{BH}_{3}+\mathrm{H}_{2}$. Isomerization of $\mathrm{BH}_{5}$ (i.e., scrambling of the nonequivalent H atoms via a $C_{2 v}$ structure) requires $9 \mathrm{kcal} / \mathrm{mol}$. Binding in $\mathrm{BH}_{5}$ is essentially due to the electron correlation; in the SCF approximation no stable $\mathrm{BH}_{5}$ is found. Although a previous CNDO calculation predicted qualitatively the equilibrium structure of $\mathrm{BH}_{3}$, it could not reveal that $\mathrm{BH}_{5}$ is hardly bound with respect to $\mathrm{BH}_{3}+\mathrm{H}_{2}$. For $\mathrm{B}_{2} \mathrm{H}_{7}{ }^{-}$the symmetric structure with a single $\mathrm{B}-\mathrm{H}-\mathrm{B}$ bond turns out to be the most stable. The calculated binding energy is $\sim-25$ $\mathrm{kcal} / \mathrm{mol}$ in SCF approximation and $\sim-28 \mathrm{kcal} / \mathrm{mol}$ with correlation, in good agreement with the experimental value of $-30 \mathrm{kcal} / \mathrm{mol}$. The change in correlation energy is less important than in other boron hydrides.

All known stable molecules with three-center two-electron bonds that involve a $\mathrm{H}_{2}$ subunit are positive ions. $\mathrm{BH}_{5}$
is the first neutral molecule of this kind; its very small binding energy results mainly from the poor electron donor properties of $\mathrm{H}_{2}$. Both in $\mathrm{BH}_{5}$ and in the "hypothetical" structures II and III of $\mathrm{B}_{2} \mathrm{H}_{7}^{-}$B is coordinated by five H atoms. Such a $\mathrm{BH}_{5}$ coordination is rather unfavorable and if it is present it is very unsymmetric, with three close and two distant H atoms.
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[^0]:    ${ }^{a}$ This value was obtained with the small basis, the corresponding values for structures $\mathrm{I}, \mathrm{II}$, and IV are $-25.6,-7.4$, and $187 \mathrm{kcal} / \mathrm{mol}$, respectively.

